# SHORTER COMMUNICATIONS

## NON-LINEAR DIFFUSION PROBLEMS WITH VARIABLE DIFFUSIVITY AND TIME-DEPENDENT FLUX BOUNDARY CONDITIONS

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- C, concentration  $\lceil \text{kg/m}^3 \rceil$ ;
- $\overline{C}_{i}, \overline{D},$ initial concentration  $\left[\text{kg}/\text{m}^3\right]$ ;
- diffusivity  $[m^2/s]$ :
- dimensionless diffusivity **;**
- & E, F, L,<br>L, L, relative change of diffusivity **:**
- function :
- coefficient :
- characteristic length [m] ;
- m exponent ;
- $\overline{T}$ , dimensionless time :
- t, time  $[s]$ ;
- $U$ , dimensionless concentration :
- x. dimensionless corodinate;
- $\mathbf{x}$ . coordinate  $[m]$ .

Greek symbols

- $\eta$ , similarity variable;<br> $\kappa$  **parameter of nonline**
- parameter of nonlinearity.

**NON-LINEAR** diffusion equations with concentrationdependent diffusion coefficient and corresponding heat conduction equations with temperature-dependent thermal properties arise in a number of physical and engineering problems. Some classes of such problems subject to the second kind boundary conditions (flux boundary conditions) have been studied by a few authors and some analytical solutions are known [1-3]. These analytical procedures, however, are valid only for the problems with constant flux boundary conditions.

The present paper establishes a general method of obtaining exact analytical solutions for a certain class of the nonlinear and variable flux type diffusion problems.

## **THEORETICAL TREATMENT**

**We are** concerned with the non-linear diffusion problem with variable diffusivity in a semi-infinite medium;

$$
\frac{\partial C}{\partial t} = \frac{\partial}{\partial x} \left[ D(C) \frac{\partial C}{\partial x} \right], \quad 0 \le x < \infty \tag{1}
$$

where  $C$  is the concentration of the diffusing substance,  $t$  the time,  $x$  the coordinate and  $D$  the diffusivity which is a positive function of the concentration. The initial and boundary conditions considered are:

$$
C = C_i \quad \text{at} \quad t = 0
$$
  

$$
D \frac{\partial C}{\partial x} = k \cdot t^{m/2} \quad \text{at} \quad x = 0
$$
 (2)

where  $k$  is an arbitrary constant and  $m$  is a positive integer or zero.

**~O~~~CLATLIR~** Let us introduce the following tr~~nsf(~rlnati(~ns\*

$$
U = \frac{D(C_i)^{m/2}}{k \cdot I^{m+1}} \int_C^{C_i} D(C) dC
$$
  
\n
$$
T = D(C_i)t/L^2
$$
  
\n
$$
X = x/L
$$
  
\n
$$
Y = D(C)/D(C_i)
$$
\n(3)

where, functional form of the dimensionless diffusivity  $\mathcal{C}(U)$ depends directly on that of the original diffusivity  $D(C)$ . In most cases of practical importance, however, it has been shown that the dimensionless diffusivity can be represented by a simple polynomial series of  $U$  and a certain parameter  $\kappa$ **[3].** When the diffusivity D(C), for instance, is described by an exponential function, the dimensionless diffusivity  $\varphi$  then becomes :

$$
\mathcal{Q}(U) = 1 - \kappa \cdot U \tag{4}
$$

where the non-linearity parameter  $\kappa$  is defined by:

$$
\kappa = \frac{k \cdot E^{n+1}}{C_i \cdot D(C_i)^{m/2+1}} \ln E
$$
  
\n
$$
E = D(C_i)/D(C=0).
$$
 (5)

Equations  $(1)$  and  $(2)$  then become:

$$
\begin{cases}\n\frac{\partial U}{\partial T} = \frac{\partial^2 U}{\partial X^2} - \kappa \cdot U \frac{\partial^2 U}{\partial X^2} \\
U = 0 \quad \text{at} \quad T = 0 \\
\frac{\partial U}{\partial X} = -T^{m/2} \quad \text{at} \quad X = 0.\n\end{cases}
$$
\n(6)

The perturbation solution for this non-linear equation is described as:

$$
U = U_1 + \kappa \cdot U_2 + \kappa^2 \cdot U_3 + \dots \tag{7}
$$

Applying the similarity analysis and the group invariant theory  $[3]$ , one can derive that:

$$
U_j(T, X) = T^{(m+1)j/2} \cdot F_j(\eta), \ j = 1, 2, 3, \dots
$$
  
\n
$$
\eta = X/2 \sqrt{T}.
$$
 (8)

Substituting equations (8) and (7) into (6) and collecting coefficients of like powers of the parameter  $\kappa$ , one can obtain the following simultaneous ordinary differential equations for the unknown similarity functions  $\vec{F}_i$  as:

$$
F''_1 + 2\eta F'_1 - 2(m+1)F_1 = 0
$$
  
\n
$$
F''_2 + 2\eta F'_2 - 4(m+1)F_2 = F_1 F''_1
$$
  
\n
$$
F''_3 + 2\eta F'_3 - 6(m+1)F_3 = F_1 F''_2 + F_2 F''_1.
$$
\n(9)

\*Where,  $L$  is a characteristic length, unit length for instance.



FIG. 1. Similarity functions  $F<sub>f</sub>$ .

The boundary conditions for these differential equations are derived from the conditions in equation (6) as:

$$
F'_1(0) = -2
$$
  
\n
$$
F'_2(0) = F'_3(0) = F'_4(0) = \dots = 0
$$
  
\n
$$
F_1(\infty) = F_2(\infty) = F_3(\infty) = \dots = 0.
$$
\n(10)

This system of the two point boundary value problems of the simultaneous linear ordinary differential equations can be easily solved by an analytical or numerical method. The dimensionless concentration  $U$ , then can be evaluated by making use of equation (8) and (7). The concentration distribution  $C(t, x)$  can be evaluated by the inverse of the transformation (3). When the attention is focused on the change of the concentration at the surface  $(x = 0)$ , we need the numerical values of  $F_j$  only at the origin ( $\eta = 0$ ).

Some examples of the similarity functions  $F_j$  and corresponding calculated results of the surface concentration changes are shown in Fig. 1, Table 1 and Fig. 2.

### **REFERENCES**

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**FIG.** 2. Change of the surface concentration.